# **Bunch Response Functions**

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### SB Amplitudes

• For the air-bag distribution  $z = \hat{z} \cos \varphi$  the offset can be expressed in terms of the SB amplitudes

$$x(\varphi,t) = \sum_{l} A_{l} \exp(il\varphi + i\chi \cos \varphi - i\Omega_{l}t) + c.c.; \quad \chi \equiv \frac{\xi \hat{z}}{R \eta}; \quad \Omega_{l} = \omega_{b} + l\omega_{s}$$

Inverse transformation follows as

$$A_{l} = \oint \frac{d\varphi}{4\pi} \left[ x(\varphi, t) + i\dot{x}(\varphi, t) / \omega_{b} \right] \exp(-il\varphi - i\chi\cos\varphi + i\Omega_{l}t);$$

A single kick distributed along the bunch as

$$\delta \dot{x} = u \cos(k_k z + \theta_k) = u \cos(q_k \cos \varphi + \theta_k)$$

gives rise to the amplitude perturbation

$$\delta A_{l} = \frac{iu}{4\omega_{h}} i^{-l} \langle k | l \rangle; \quad \langle k | l \rangle = e^{i\theta_{k}} J_{l}(\chi - q_{k}) + e^{-i\theta_{k}} J_{l}(\chi + q_{k})$$

where the bracket factor  $\langle k | l \rangle$  describes the kicker visibility for the mode l.

#### Offset distribution

In case when there is an infinite number of kicks applied with a given frequency as  $u_n = V \cos(\omega t_n + \psi)$ ;  $t_n = nT_0$ , they are summarized giving

$$\delta A_{l} = \frac{iV}{4\omega_{b}} i^{-l} \langle k | l \rangle \frac{1}{2} \frac{\exp(i(\Omega_{l} - \omega)t_{n} - i\psi)}{1 - \exp(-i(\Omega_{l} - \omega)T_{0})}$$

Taking into account that only small interval of frequencies around the eigen-frequency  $\Omega_i$  is of interest, the exponent in the denominator is expanded leading to

$$x(\varphi,t) = \operatorname{Re}\left\{\frac{Ve^{-i(\omega t + \psi)}}{4\omega_b}i^{-l}\sum_{l}\langle k|l\rangle\frac{\exp(il\varphi + i\chi\cos\varphi)}{\Omega_l - \omega}\right\}$$

where all the eigen-frequencies lie in the lower half-plane.

### Measured Response

Let the pickup measurement be modeled as

$$R = \oint \frac{d\varphi}{4\pi} \left( x(\varphi) + x(-\varphi) \right) \cos(q_p \cos \varphi + \theta_p)$$

Integration results in

$$R(t) = \frac{V\sqrt{\beta_k \beta_p}}{8C} \operatorname{Re} \left[ e^{-i(\omega t + \psi)} G(\omega) \right]; \quad G(\omega) = \sum_{l} \frac{\langle k | l \rangle \langle l | p \rangle}{\Omega_l - \omega},$$

with the bracket

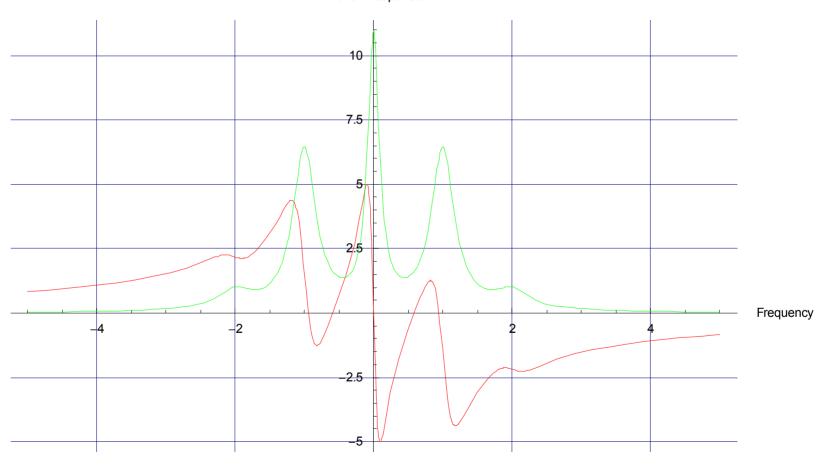
$$\langle l | p \rangle = e^{-i\theta_p} J_l(\chi - q_p) + e^{i\theta_p} J_l(\chi + q_p)$$

describing how effective is the mode seen by the pickup.

## Air-bag Response

• Air-bag response functions (Re and Im parts) for  $N/N_{th} << 1$  and  ${\rm Im}\,\Omega_{_{l}} = (0.1\,|\,l\,|\,+0.1)\omega_{_{s}}$ 





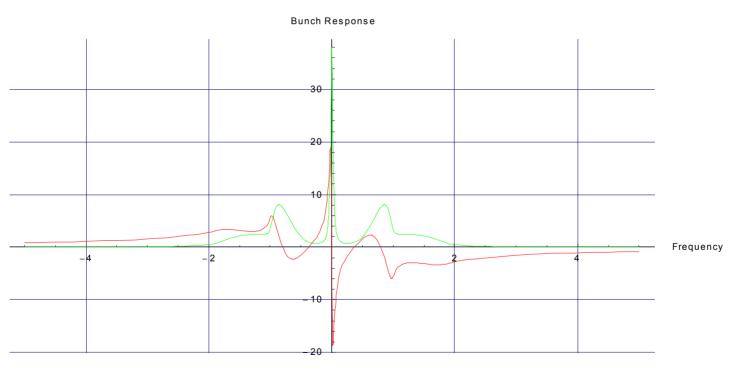
## Parabolic distribution with non-linearity

A model distribution

$$f(\hat{z}), \,\omega_s(\hat{z}) \propto 1 - \hat{z}^2/z_m^2$$

response is calculated by integration of it's air-bags with

$$\Omega_l = l\omega_s(\hat{z}) - 0.02i \omega_s(0)$$



#### If not far from the threshold

- When the bunch population is not far from the strong headtail threshold (TMCI), the impedance modifies the response:
  - > They are not symmetric any more
  - > The peaks are shifted
  - > The peaks are not equidistant
- Although the response can be formally expressed in the same way,

$$R(t) = \frac{V\sqrt{\beta_k \beta_p}}{8C} \operatorname{Re} \left[ e^{-i(\omega t + \psi)} G(\omega) \right]; \quad G(\omega) = \sum_{l} \frac{\langle k | l \rangle \langle l | p \rangle}{\Omega_l - \omega},$$

the visibility form-factors are not as easy to calculate.